

# Solving the radiation equation

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Heat can transfer at the speed of light through radiation. The **Stefan-Boltzmann Law** describes how the power  $E$  emitted from an object depends on its temperature:

$$E = \varepsilon\sigma T^4 \quad (1)$$

where  $\sigma$  is called the Stefan constant,  $\varepsilon$  is the emissivity (1 if the object is a perfect black body), and  $T$  is the temperature in Kelvin.

For a system that has many objects, we have to discretize every one of them into many small patches  $\{dA_i\}$  in order to model their interactions. The power density given away by each patch is called **radiosity** (denoted by  $B$ ), which is governed by the following integral equation (this is in fact an inhomogeneous Fredholm equation of the second kind<sup>1</sup>):

$$B_i dA_i = E_i dA_i + \rho_i \int_{\Omega} B_j F_{ji} V_{ji} dA_j \quad (2)$$

where  $B_i$  is the outgoing power per unit area from patch  $i$ ,  $E_i$  is the emission power per unit area of the patch,  $\rho_i$  is the reflectivity of the patch,  $dA_i$  is the area of the patch,  $F_{ji}$  is the view factor (dimensionless), and  $V_{ji}$  is the visibility function (0 if patches  $i$  and  $j$  cannot see each other, otherwise 1). The visibility function can be calculated using collision detection between rays and patches.

The view factor from a surface  $A_1$  to a surface  $A_2$  defines the proportion of the radiation leaving  $A_1$  that ends up striking  $A_2$ . In 3D, it can be calculated through the following formula based on the etendue of a ray:

$$F_{12} = \frac{1}{A_1} \int_{A_1} \int_{A_2} \frac{\cos \theta_1 \cos \theta_2}{\pi r_{12}^2} dA_1 dA_2 \quad (3)$$

where  $r_{12}$  is the distance between the two surfaces, and  $\theta_1$  and  $\theta_2$  are the acute angles between the surface normal and the distance vector. In 2D, it can be calculated by replacing the areas with lengths and  $r_{12}^2$  with  $r_{12}$ :

$$F_{12} = \frac{1}{L_1} \int_{L_1} \int_{L_2} \frac{\cos \theta_1 \cos \theta_2}{\pi r_{12}} dL_1 dL_2 \quad (3a)$$

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<sup>1</sup> [http://en.wikipedia.org/wiki/Fredholm\\_integral\\_equation](http://en.wikipedia.org/wiki/Fredholm_integral_equation)

The net energy that converts into thermal energy at a patch comes from the difference between the portion of energy reflected from all other patches that becomes absorbed by the patch and the energy that it emits:

$$Q_i dA_i = \alpha_i \int_{\Omega} B_j F_{ji} V_{ji} dA_j - E_i dA_i \quad (4)$$

where  $\alpha_i$  is the absorptivity (the ability to absorb electromagnetic radiation) of the patch's material. This equation can be coupled with the Heat Equation and the Navier-Stokes Equation to allow radiation to interact with the flow of heat and mass.

Equations (2) and (4) can be discretized into the following matrix equations:

$$B_i - \rho_i \sum_j F_{ij} V_{ij} B_j = E_i \quad (5)$$

$$Q_i = \alpha_i \sum_j F_{ij} V_{ij} B_j - E_i \quad (6)$$

Equation (5) can be solved using a relaxation method, such as the Gauss-Seidel method,<sup>2</sup> using the emission power at each patch as the initial guess. In a transient simulation, only a few relaxation steps are needed as the time-dependent solution is already somehow iterative.

The thermal power generated at each patch  $Q_i$  can then be derived using Equation (6).

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<sup>2</sup> [http://en.wikipedia.org/wiki/Gauss%E2%80%93Seidel\\_method](http://en.wikipedia.org/wiki/Gauss%E2%80%93Seidel_method)